

Burgers equation and Fourier Fractal Decimation



Bangalore, ICTS
24 October 2014

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ERC Advanced Grant (N. 339032) "NewTURB"
(P.I. Prof. Luca Biferale)



European Research Council
Established by the European Commission

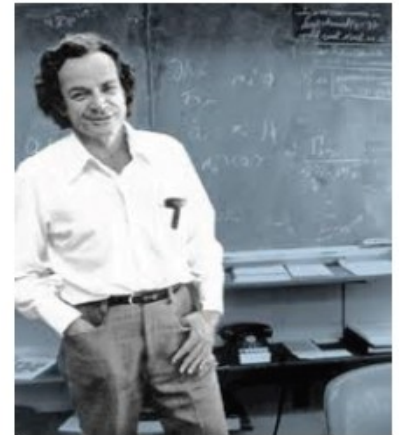
Outline:

1) Why are **Navier-Stokes** equations interesting for Theoretical Physics?

- Strongly non perturbative field Theory (Classical)
- Anomalous Scaling (Non-Gaussian Statistics)

2) Why do we need a model for Navier-Stokes?

3) Burgers' equation and Fourier Fractal Decimation



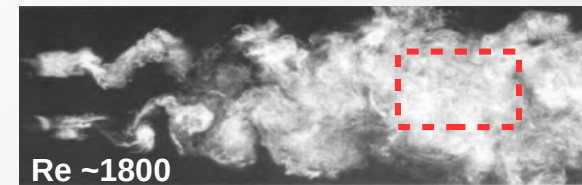
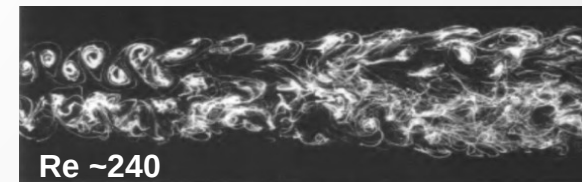
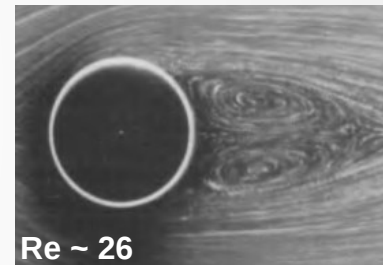
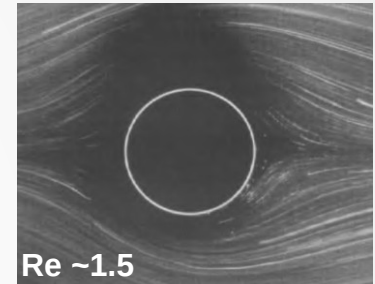
“With turbulence, it's not just a case of physical theory being able to handle only simple cases—we can't do any. We have no good fundamental theory at all.” (Feynman, 1979, Omni Magazine, Vol. 1, No.8).

Probabilistic description for fully developed Turbulence

Navier-Stokes, (N-S), equations:

$$\begin{cases} \frac{\partial \mathbf{v}(\mathbf{x}, t)}{\partial t} + \mathbf{v}(\mathbf{x}, t) \cdot \nabla_x \mathbf{v}(\mathbf{x}, t) = -\nabla_x p(\mathbf{x}, t) + \nu \Delta_x \mathbf{v}(\mathbf{x}, t) + \mathbf{f}(\mathbf{x}, t) \\ \nabla_x \cdot \mathbf{v}(\mathbf{x}, t) = 0 \end{cases}$$

$$\begin{cases} \hat{t} = t/t_0 \\ \hat{x} = x/l_0 \\ \hat{v} = v/v_0 \end{cases} \quad \partial_t \hat{v} + \hat{v} \cdot \partial \hat{v} = -\partial \hat{P} + \frac{1}{Re} \partial^2 \hat{v} \quad Re = \frac{l_0 v_0}{\nu} \quad Re \sim \frac{\hat{v} \partial \hat{v}}{\nu \partial^2 \hat{v}}$$



Left-right invariance is broken

Z-invariance is broken, discrete time invariance

At high Re symmetries are spontaneously broken

Restored symmetries; (in a statistical sense)

$Re \sim 10$

$\sim 10^2$

$\sim 10^3$

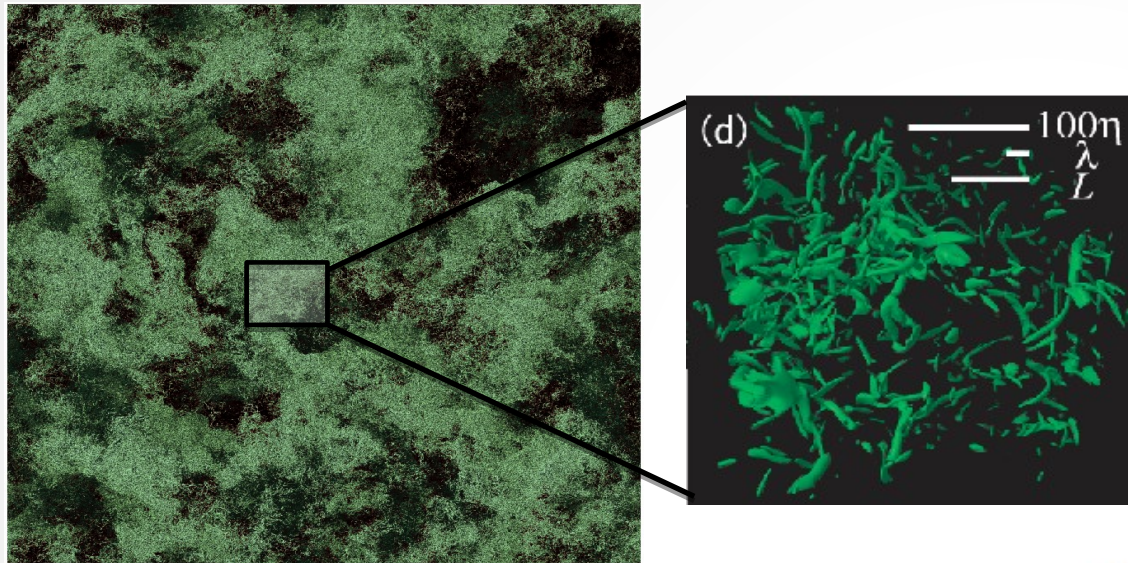
Recirculating standing eddies

Kàrmàn street

Flow becomes chaotic in its time-dependence

Homogeneous-isotropic fully developed turbulence

Anomalous Exponents, Small-Scales Intermittency

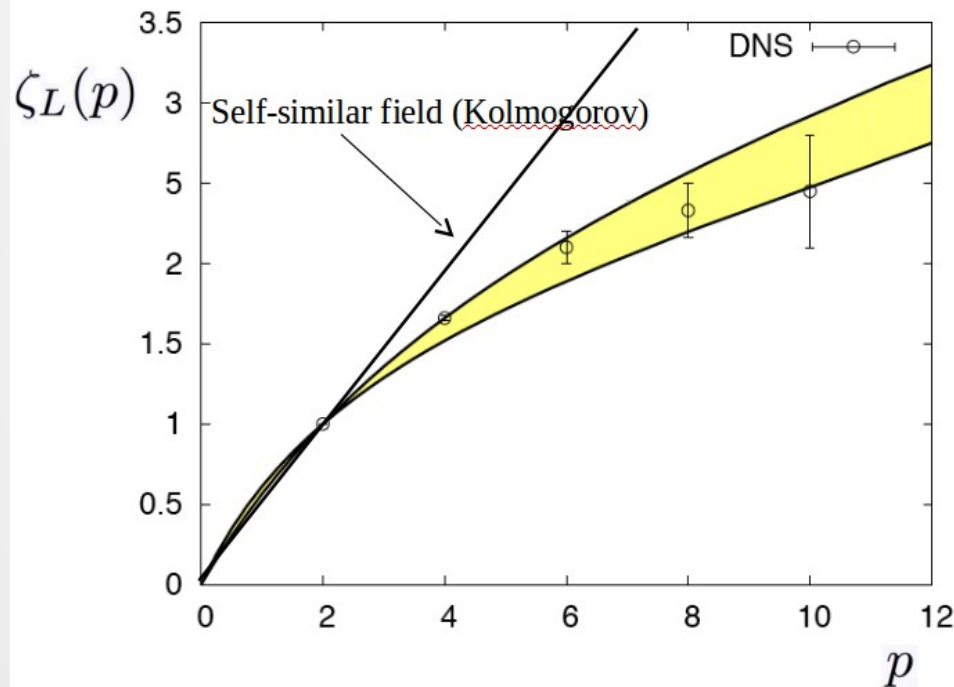


H1) Restored symmetries (in a statistical sense).

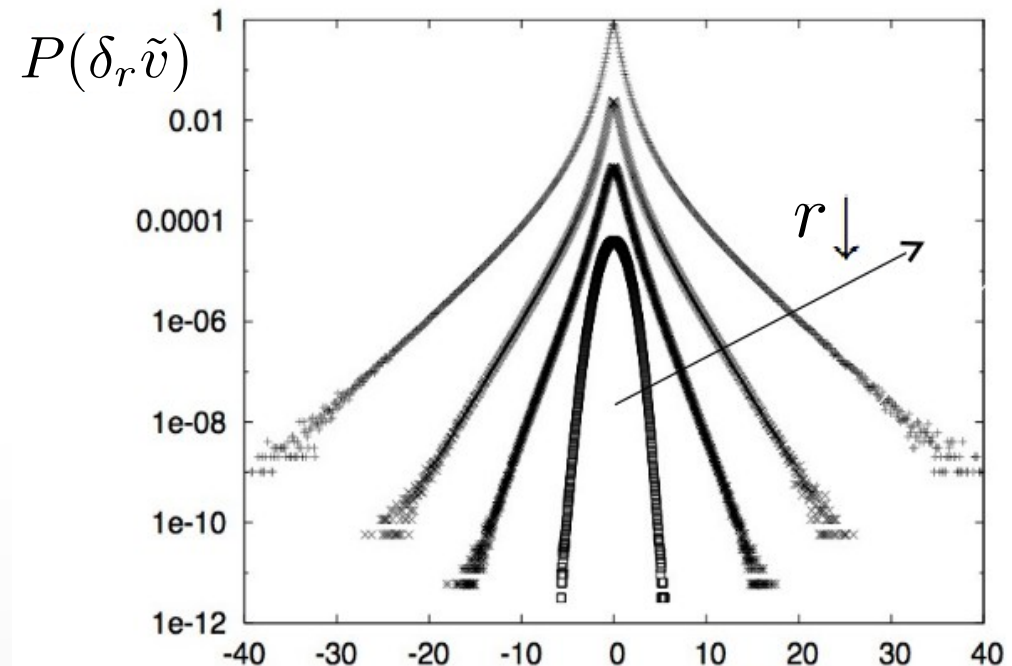
H2) Self-similarity at small scales.

$$S_p(r) = \langle (\delta_r v)^p \rangle \sim r^{\zeta(p)}$$

$$\delta_r v = v(x+r) - v(x)$$



Benzi, Biferale, Fisher, Lamb and Toschi, JFM (2010).



$$\delta_r \tilde{v} = \frac{\delta_r v}{\langle (\delta_r v)^2 \rangle^{1/2}}$$

..a model for Turbulence

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2}$$

Burgers' equation

$u(t, x)$: velocity field, depending on a variable of time (t), and on a variable of space (x) | ν : kinematic viscosity

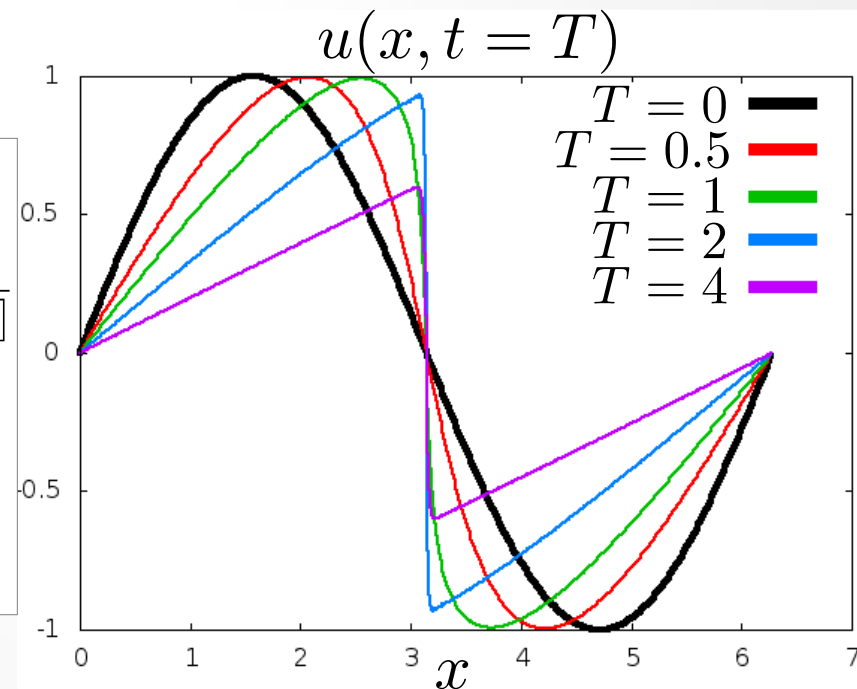
Burgers produces a singularity, (shock).

Lagrangian observation

$$\begin{cases} u(t, X(t, a)) = u_0(a) \\ X(t, a) = a + tu_0(a) \end{cases}; J(t, a) = \frac{\partial X}{\partial a} = 1 + tu'_0(a); t^* = \frac{1}{-\text{inf}_a[u'_0(a)]}$$

Gradient in the Eulerian coordinates

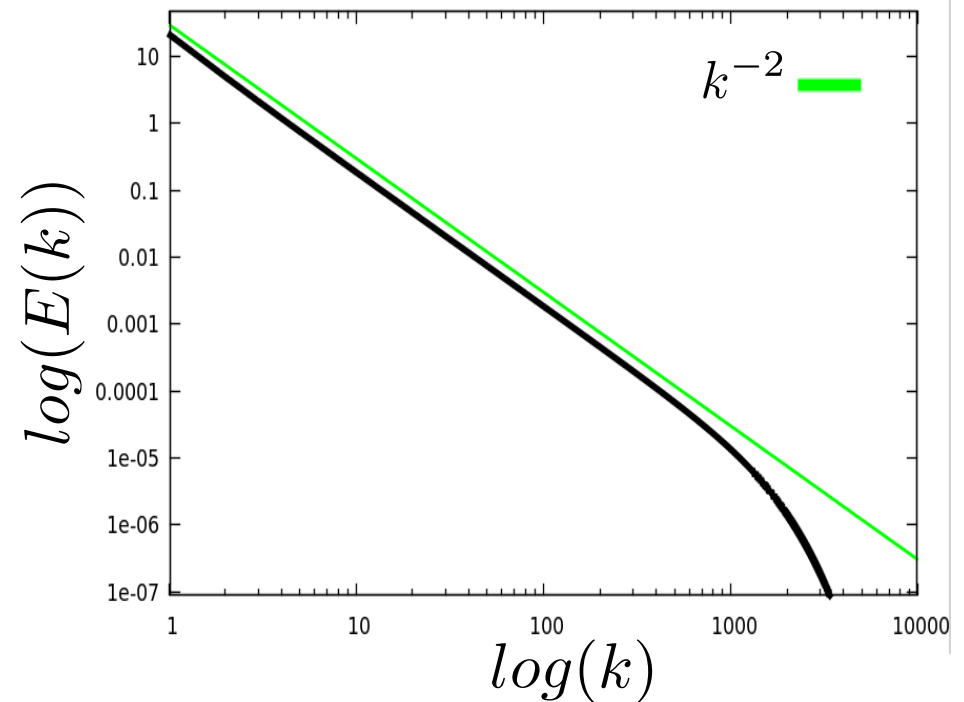
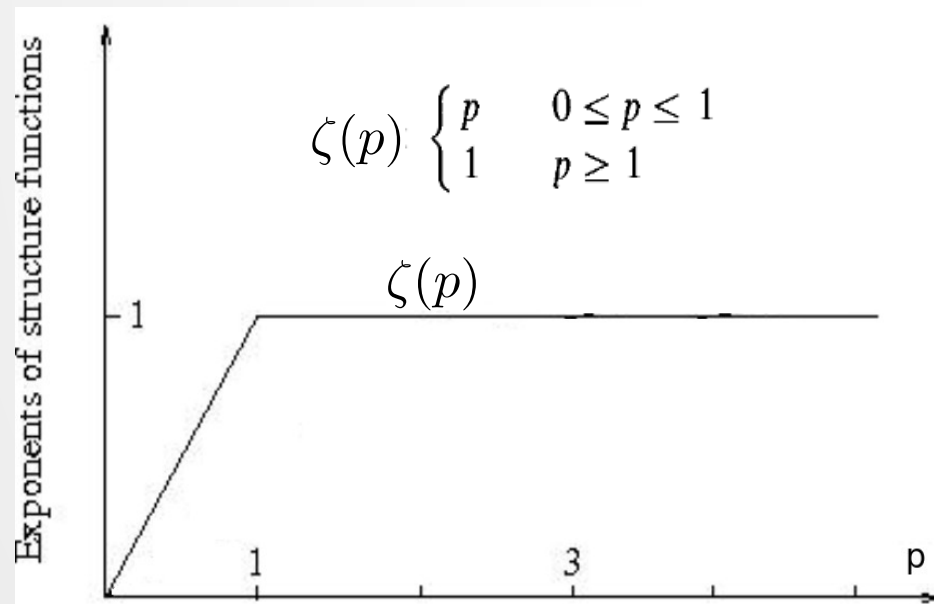
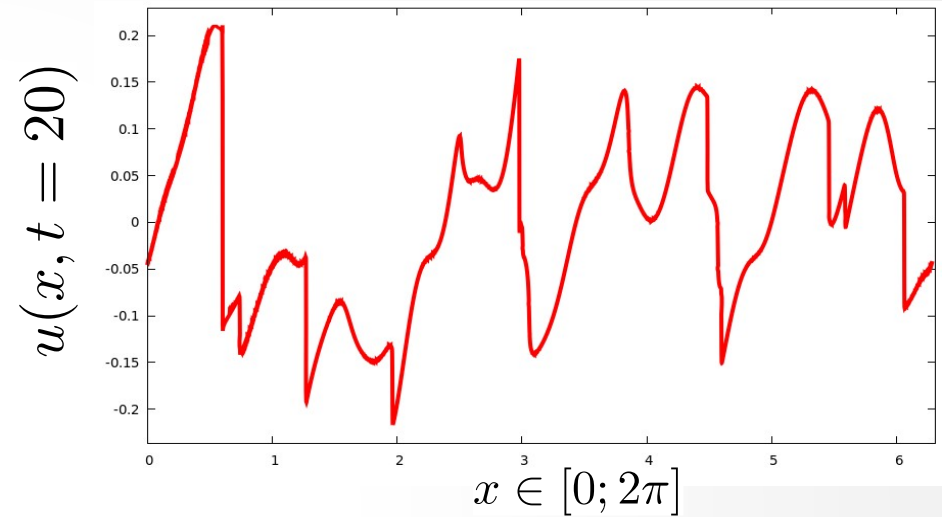
$$\frac{\partial u}{\partial x} \Big|_{x^*=a^*} = \frac{\partial u}{\partial a} \Big|_{a^*} \frac{\partial a}{\partial x} \Big|_{x^*} = u'_0(a) \frac{1}{1 + tu'_0(a)} \rightarrow \lim_{t \rightarrow t^*} \frac{u'_0(a)}{1 + tu'_0(a)} = \infty$$



Intermittency on Burgers' equation

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2} + f(x, t)$$

$$S_p(r) = \langle (\delta_r v)^p \rangle \sim r^{\zeta(p)}$$



how many degrees of freedom are related to the singularity?

..Reduce to learn!

FRACTAL FOURIER DECIMATION

$$u(x, t) = \sum_{k \in \mathbb{Z}} e^{ikx} u(k, t) \quad P_D \cdot u(x, t) = \sum_{k \in \mathbb{Z}} e^{ikx} \theta_k u(k, t)$$

$$\theta_k = \begin{cases} 1 & \text{with probability } h_k \\ 0 & \text{with probability } 1 - h_k, \quad k \equiv |\mathbf{k}| \end{cases}$$

$$h_k = (k/k_0)^{D-1}, \quad 0 < D \leq 1$$

The decimation is Random but Quenched on time,

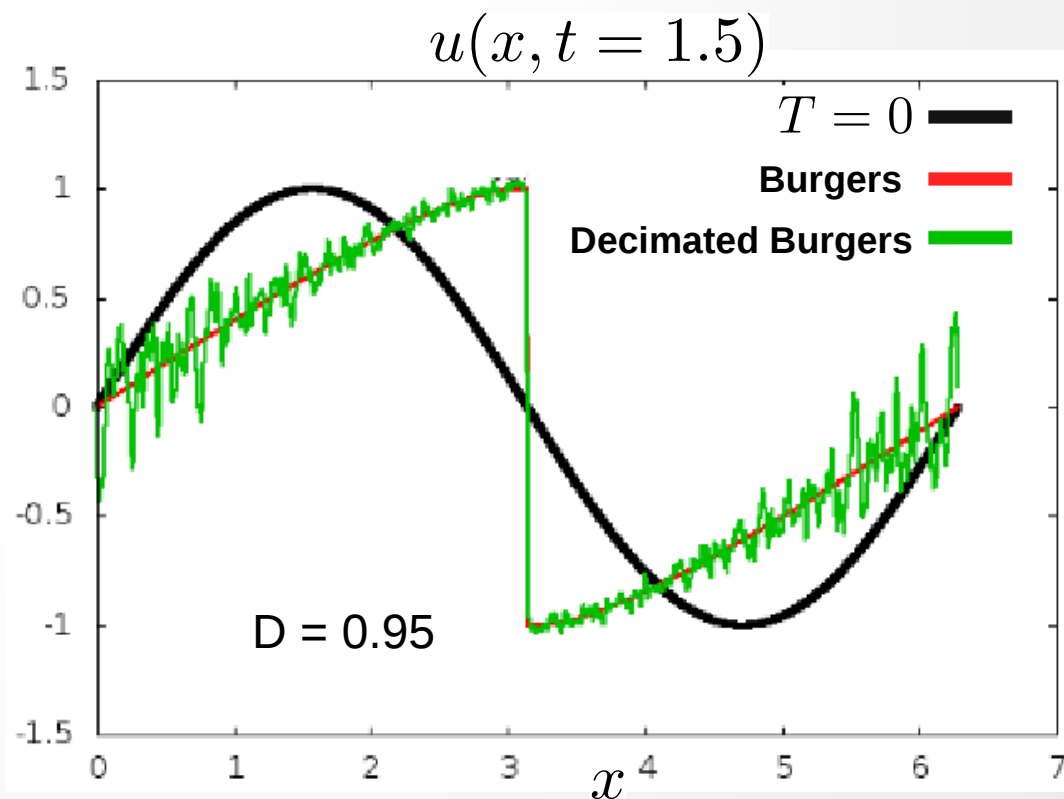
leaving on average $N(k) \sim k^D$ active mode

Galerkin truncation projection: $k < k_{max}$



- Finite number of d.o.f.
- Fractal dimension

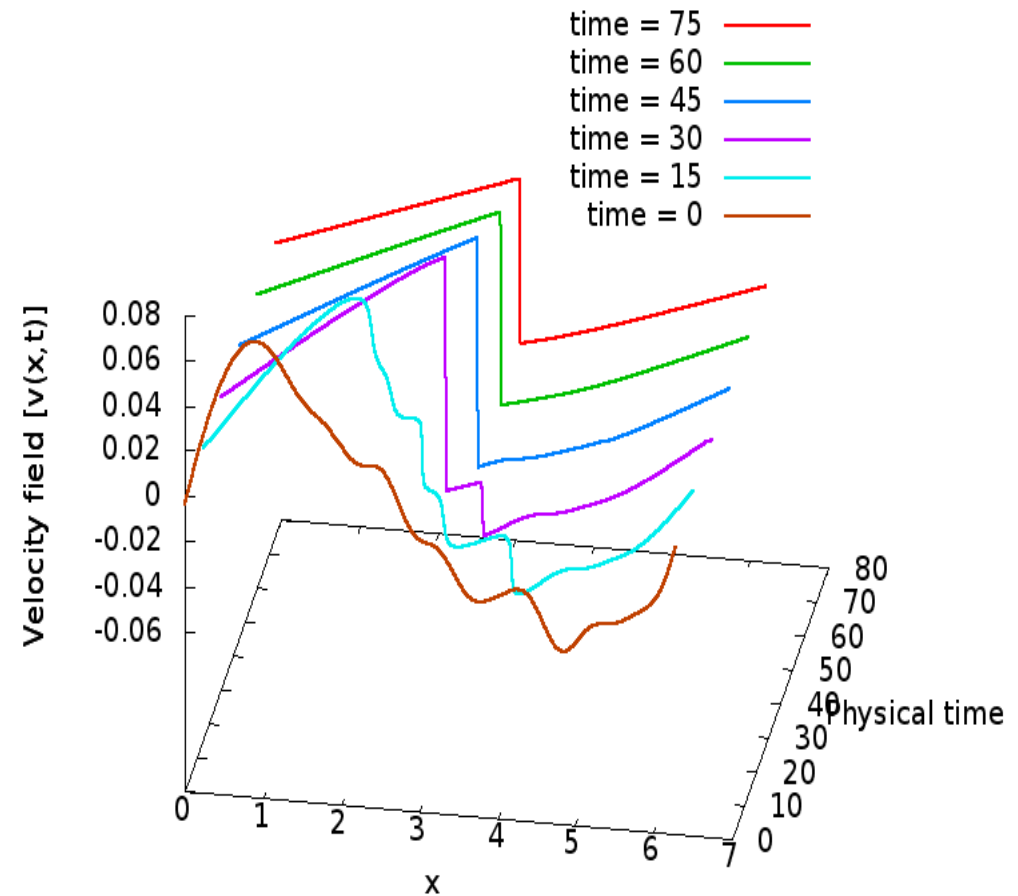
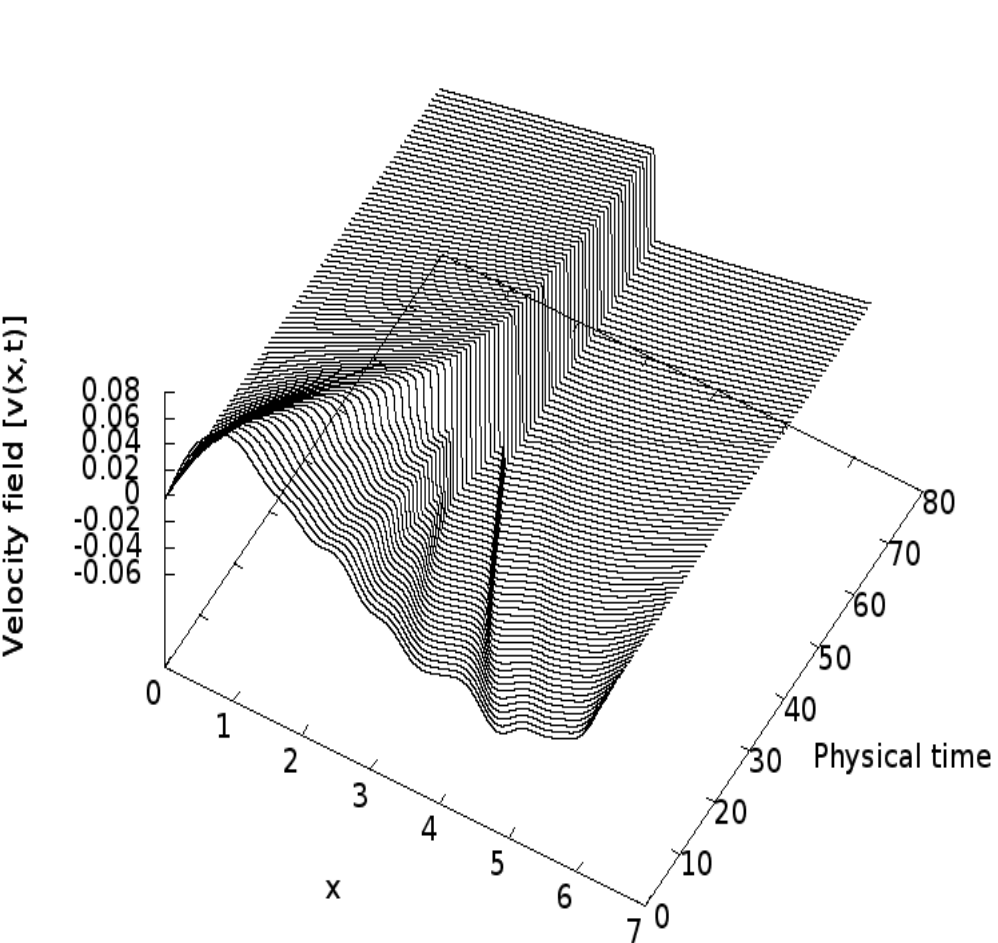
Frisch, Pomyalov, Procaccia, and Ray,
Turbulence in non-integer dimensions by
fractal Fourier decimation. Phys. Rev. Lett.
108, (2012)



Real space velocity field evolution:

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2}$$

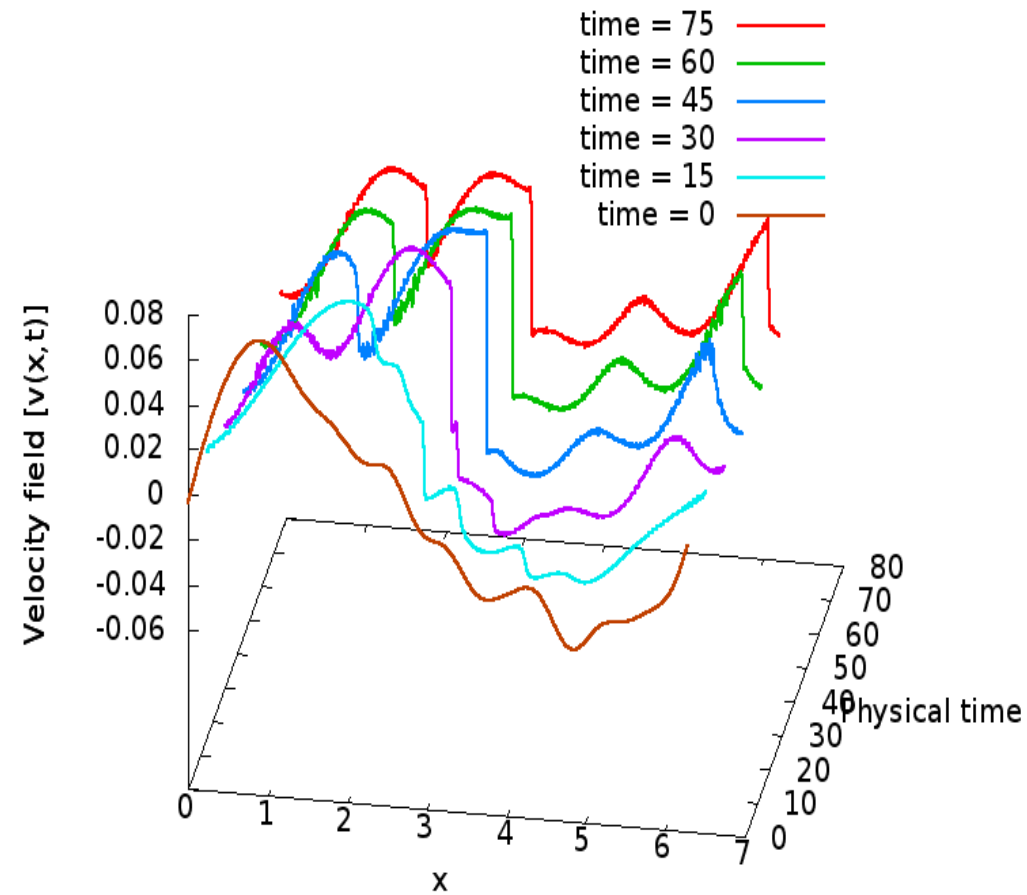
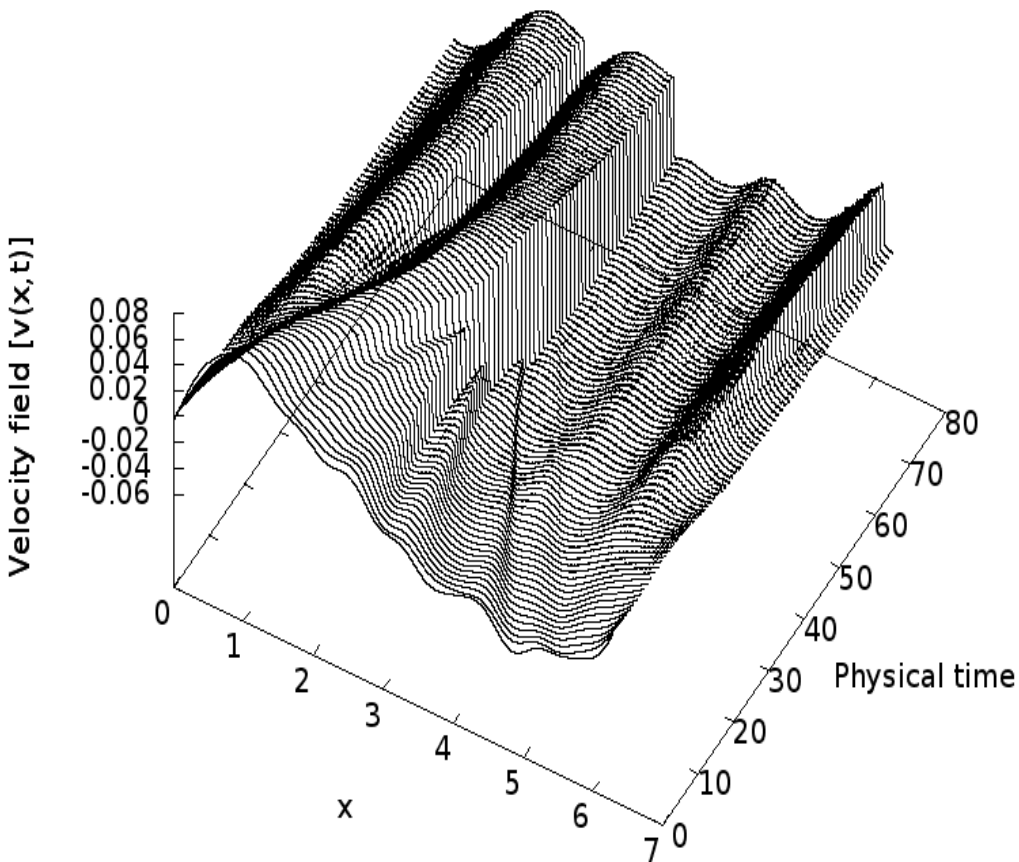
$Dim = 1$



Real space velocity field evolution:

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2}$$

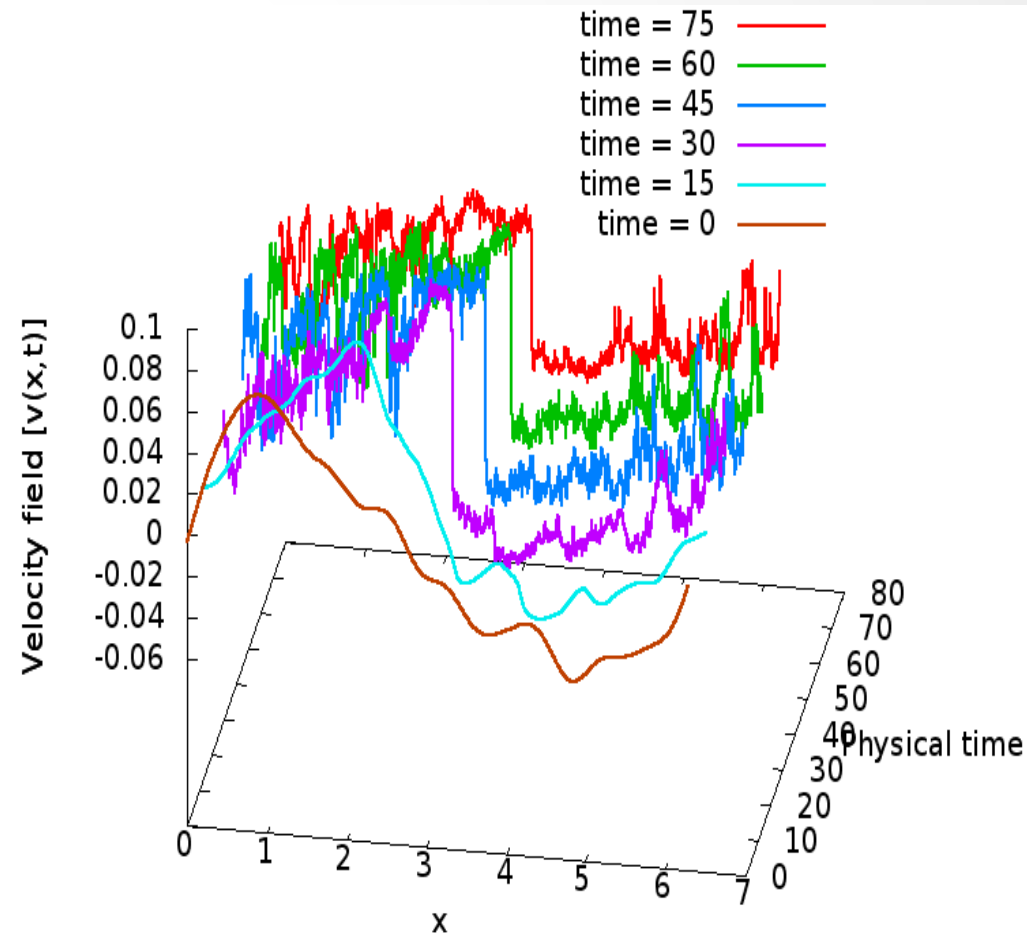
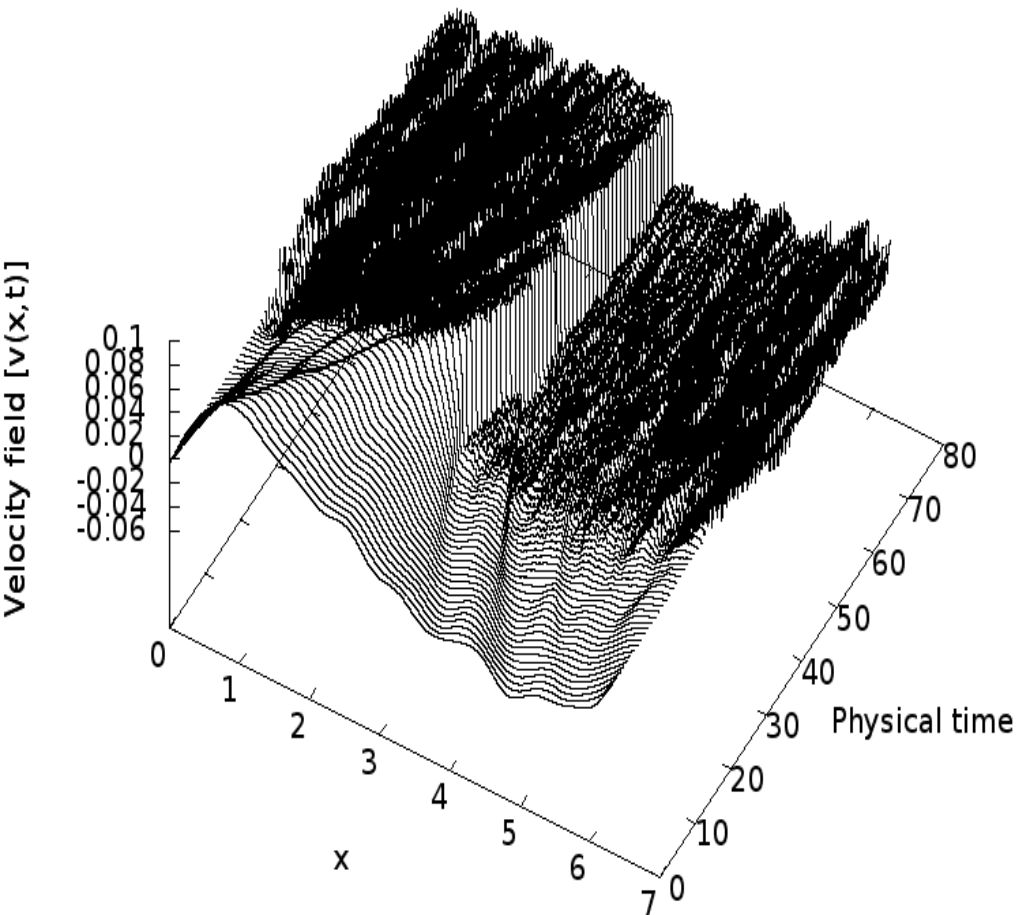
$$Dim = 0.99$$



Real space velocity field evolution:

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2}$$

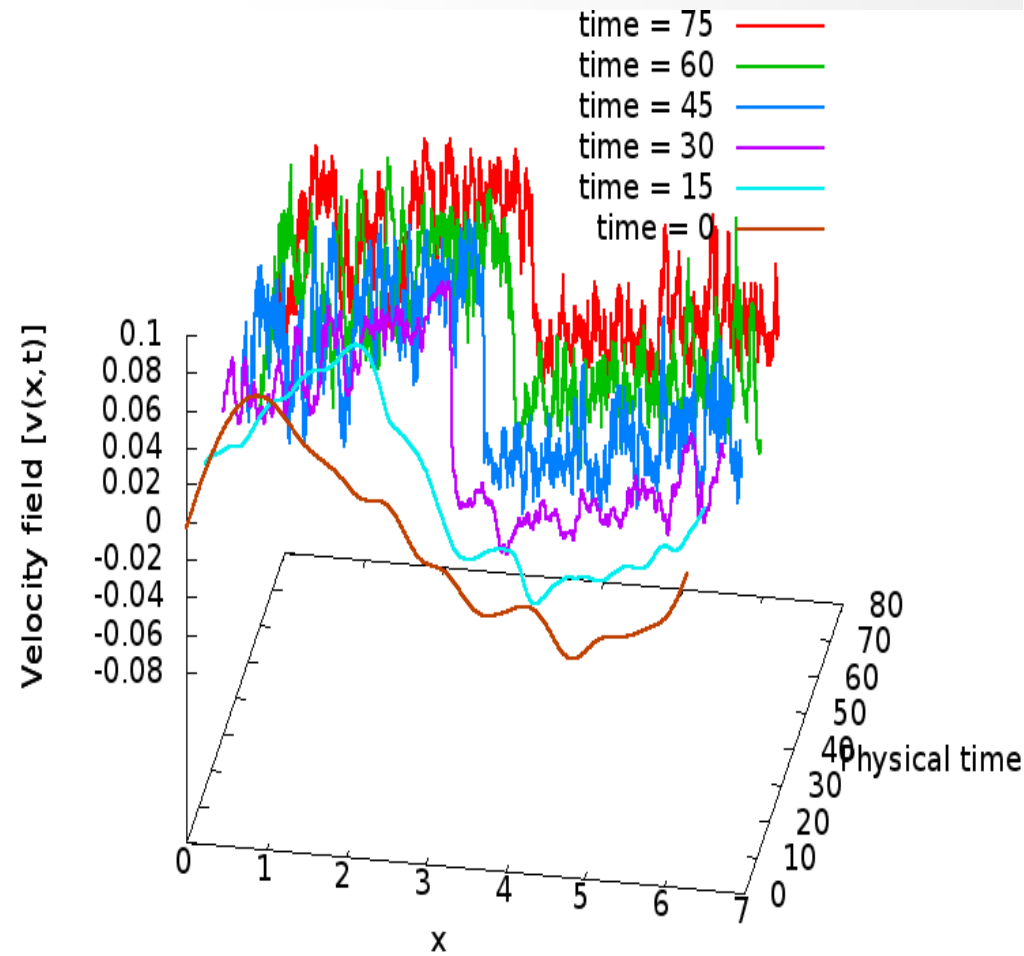
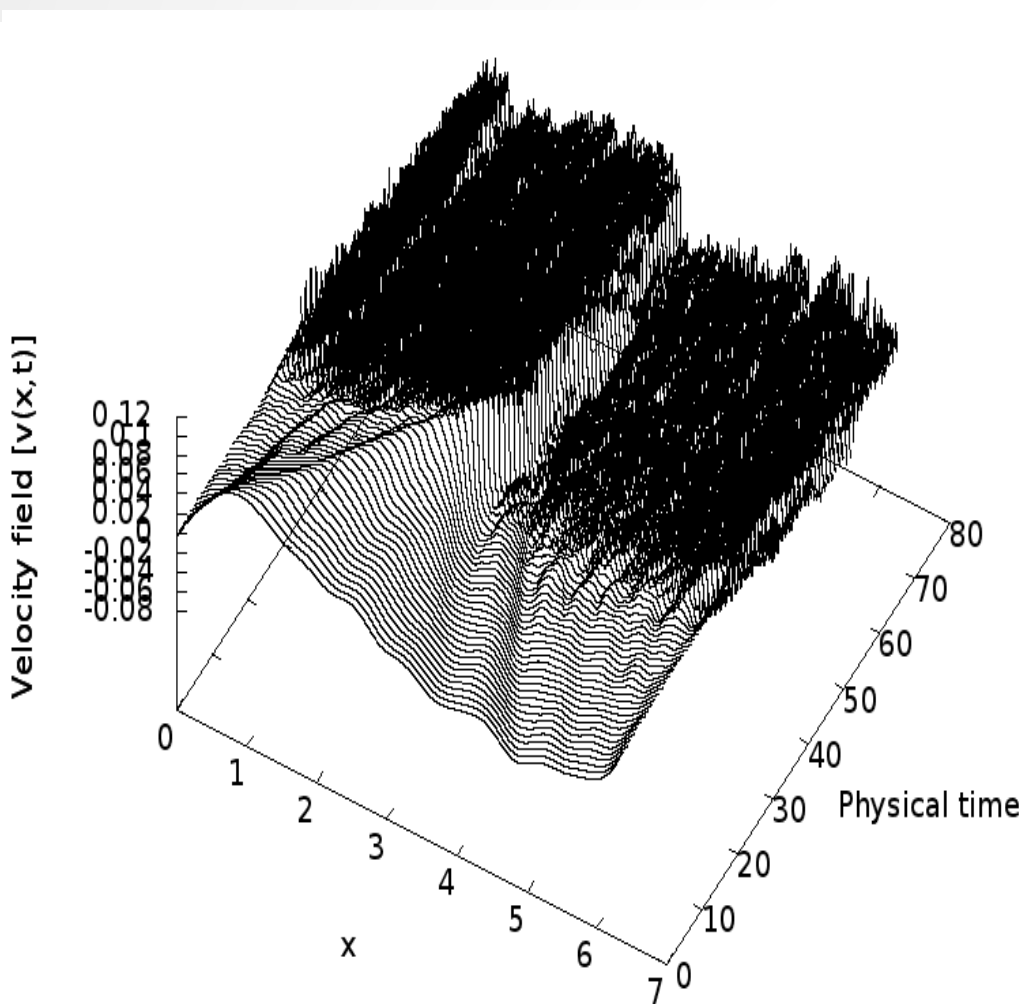
$$Dim = 0.95$$



Real space velocity field evolution:

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2}$$

$$Dim = 0.85$$



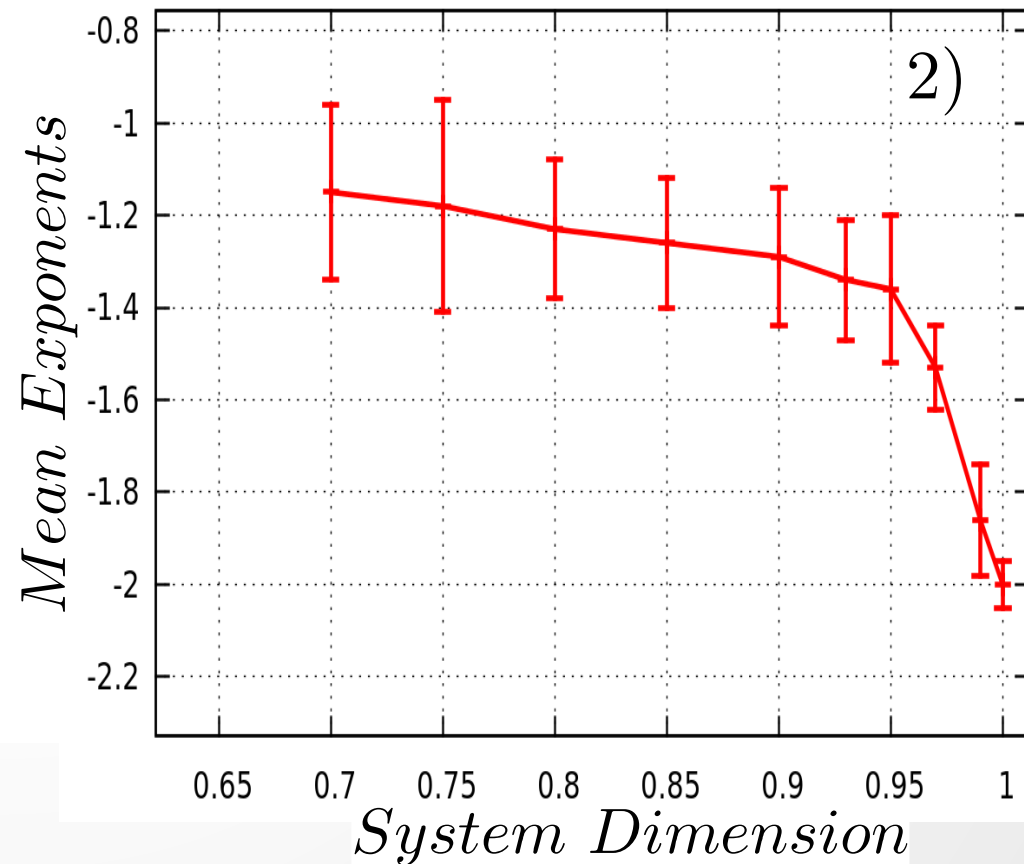
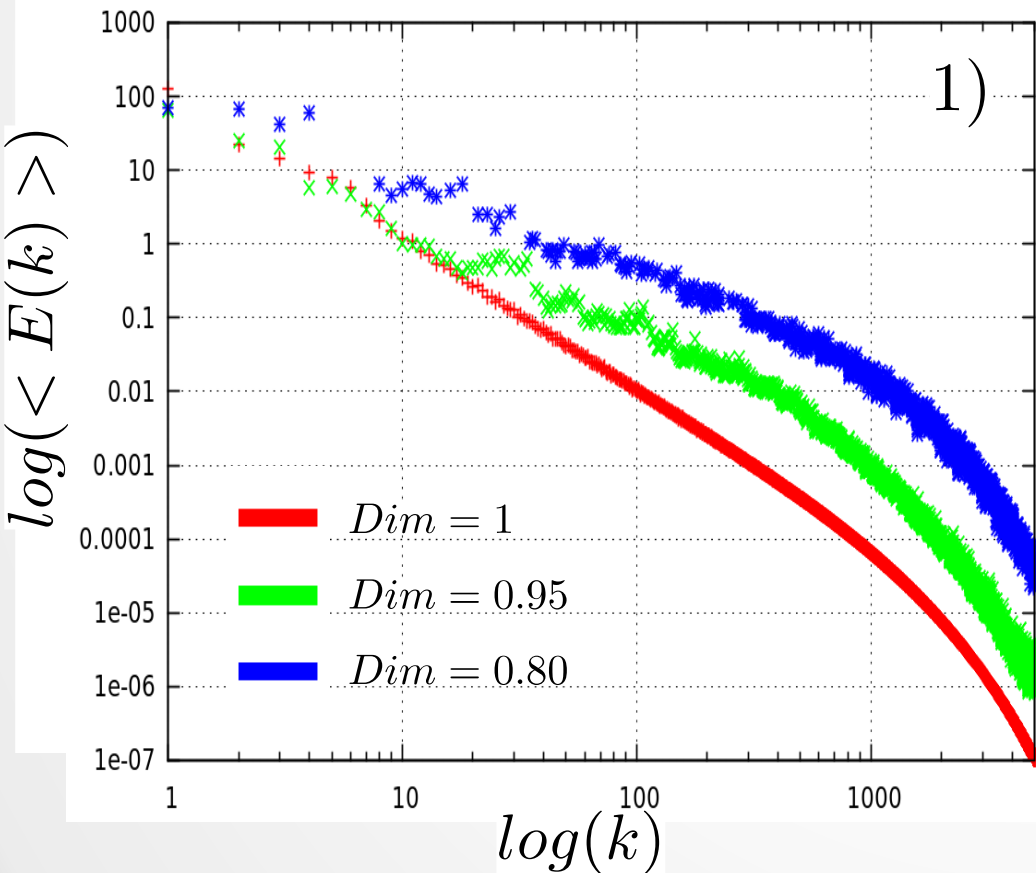
Decimated Energy Spectrum

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2} + f(x, t)$$

Mean spectra:

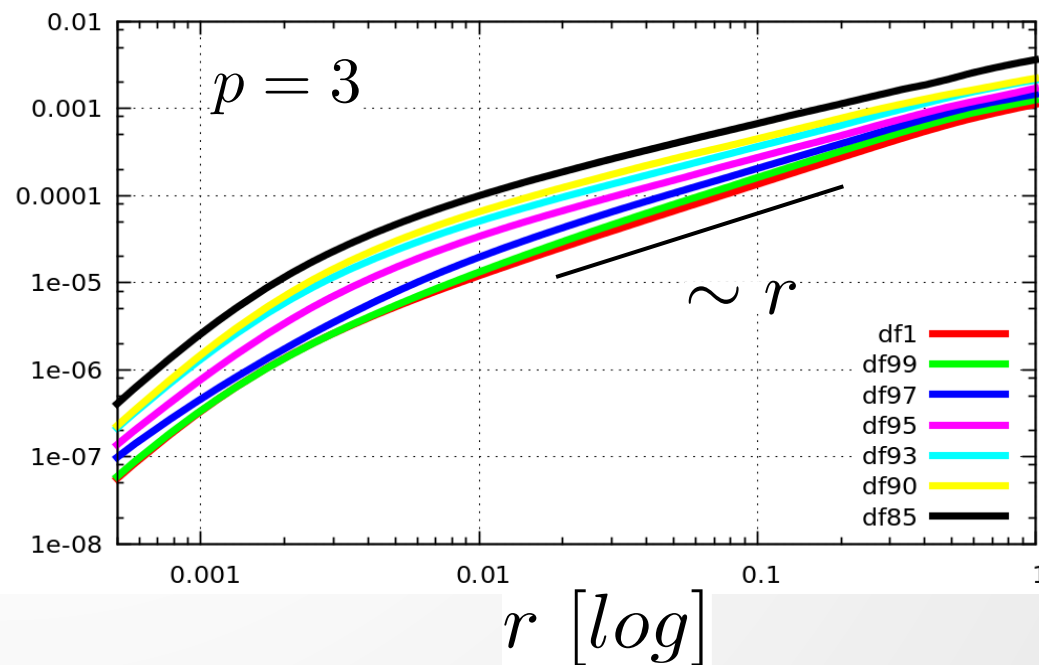
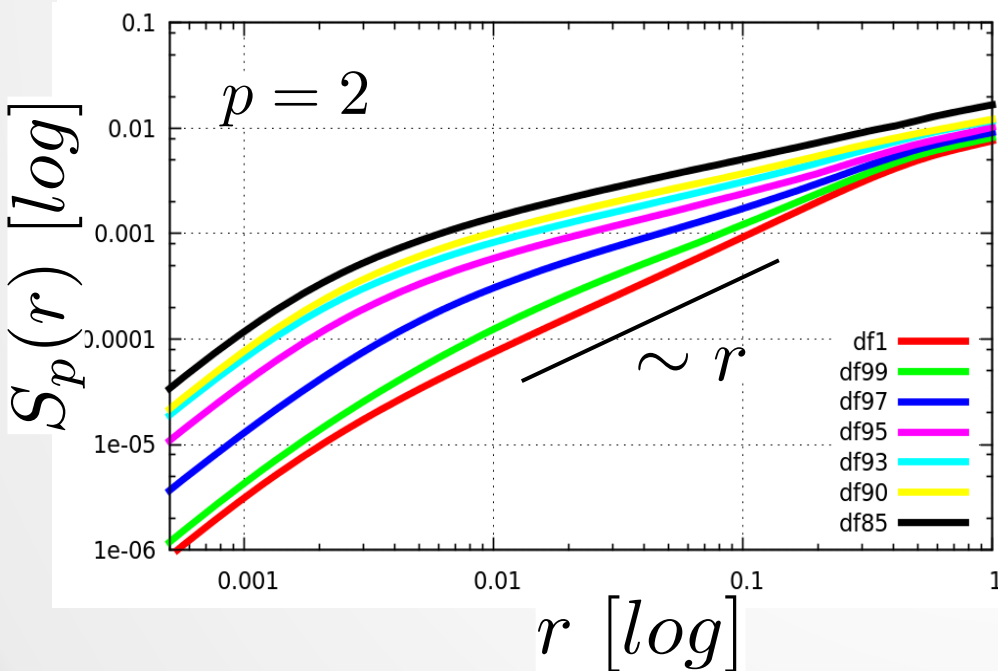
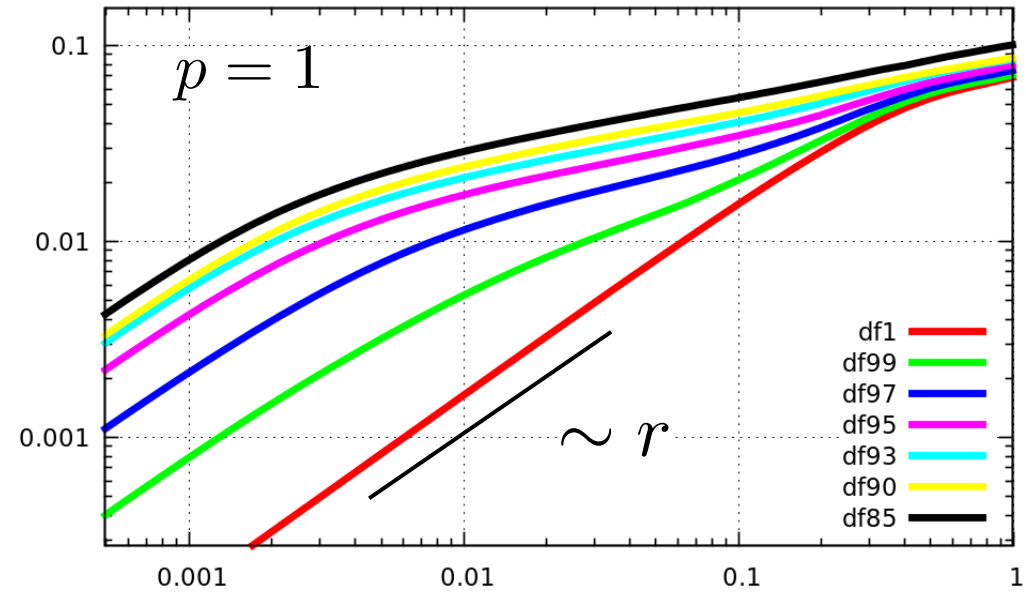
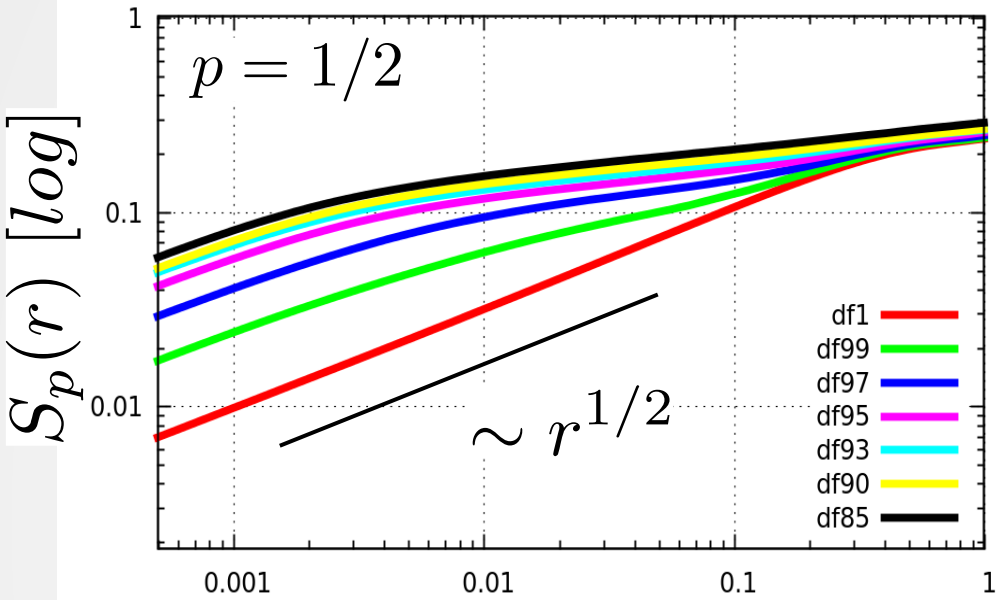
1) Single Mask: Averaged on time in the stationary state

2) Mean Slope (32 masks) vs Fractal Dimension



Structure Functions

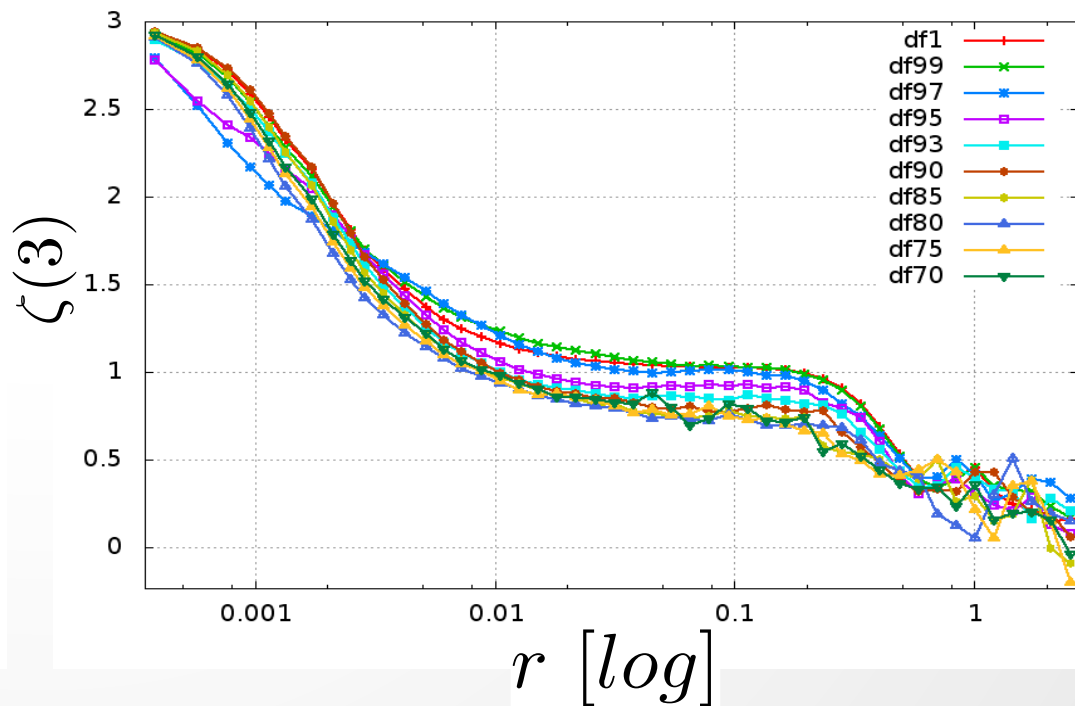
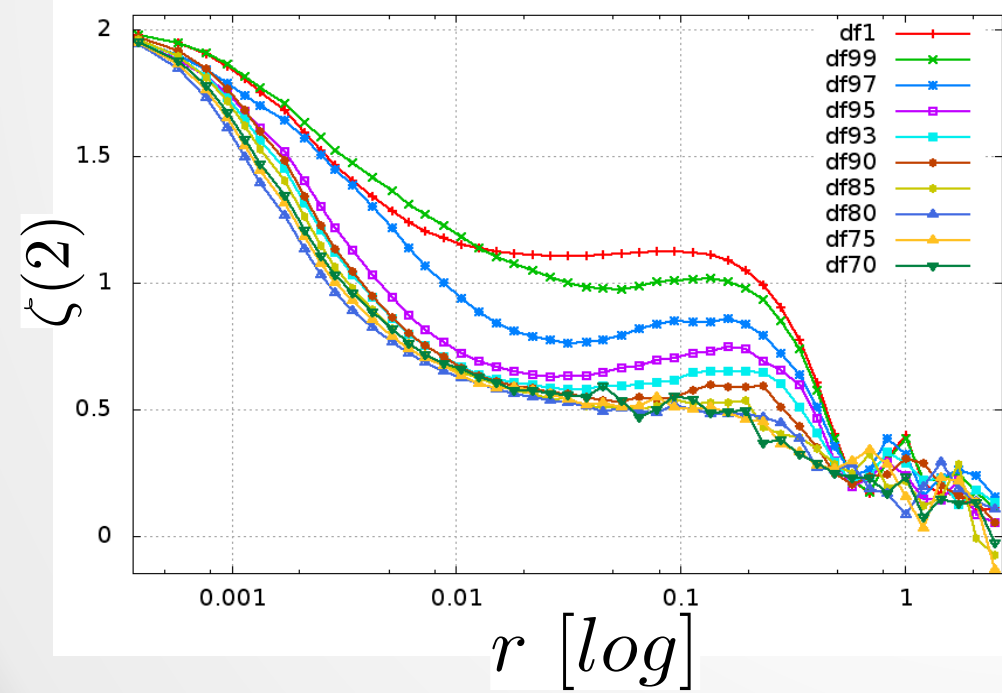
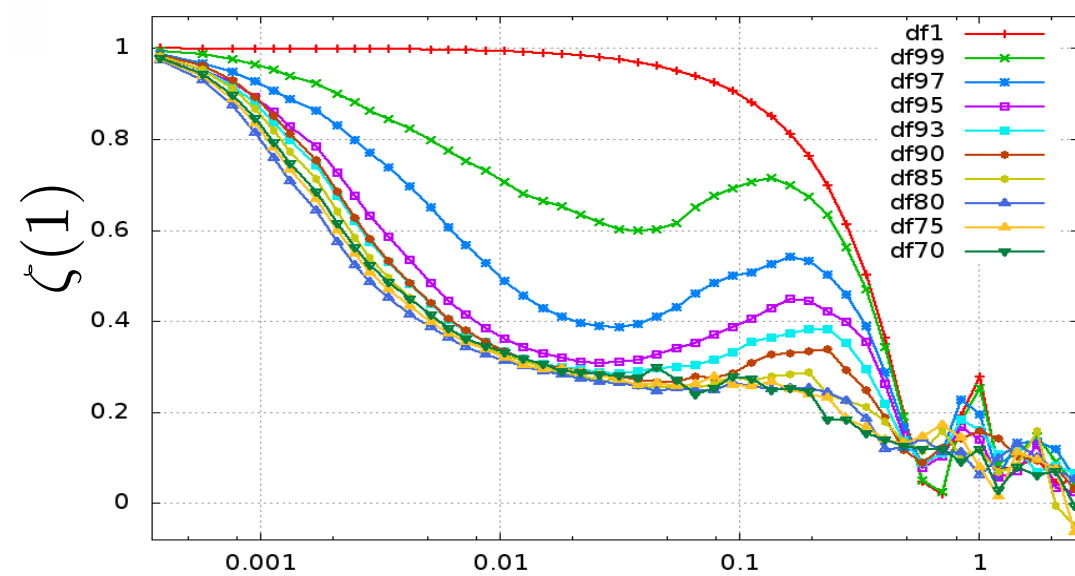
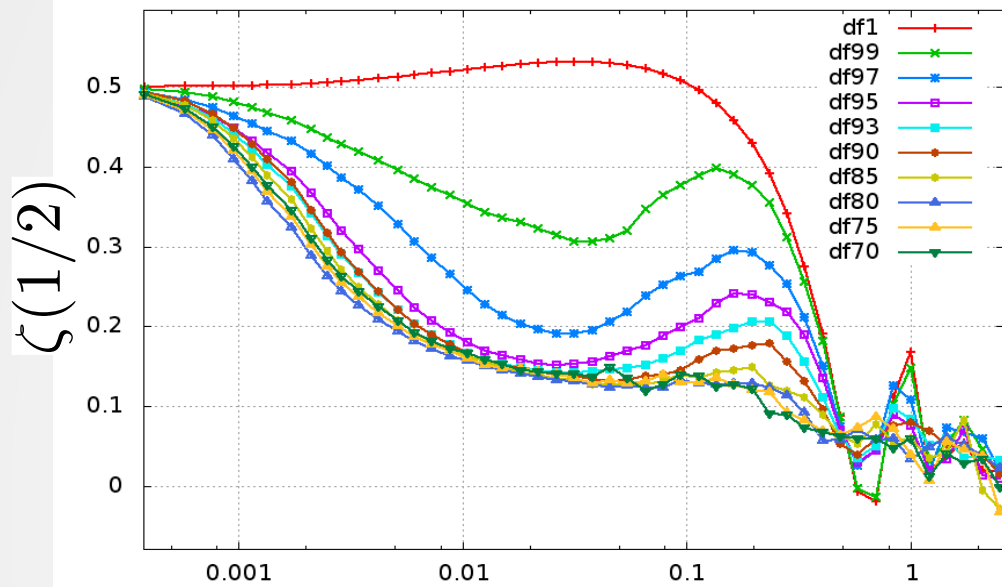
$$S_p(r) = \langle (\delta_r v)^p \rangle \sim r^{\zeta(p)}$$



Structure Functions: Local Slope

$$S_p(r) = \langle (\delta_r v)^p \rangle \sim r^{\zeta(p)}$$

$$\zeta(p) = \frac{\partial \log(S_r(r))}{\partial \log(r)}$$

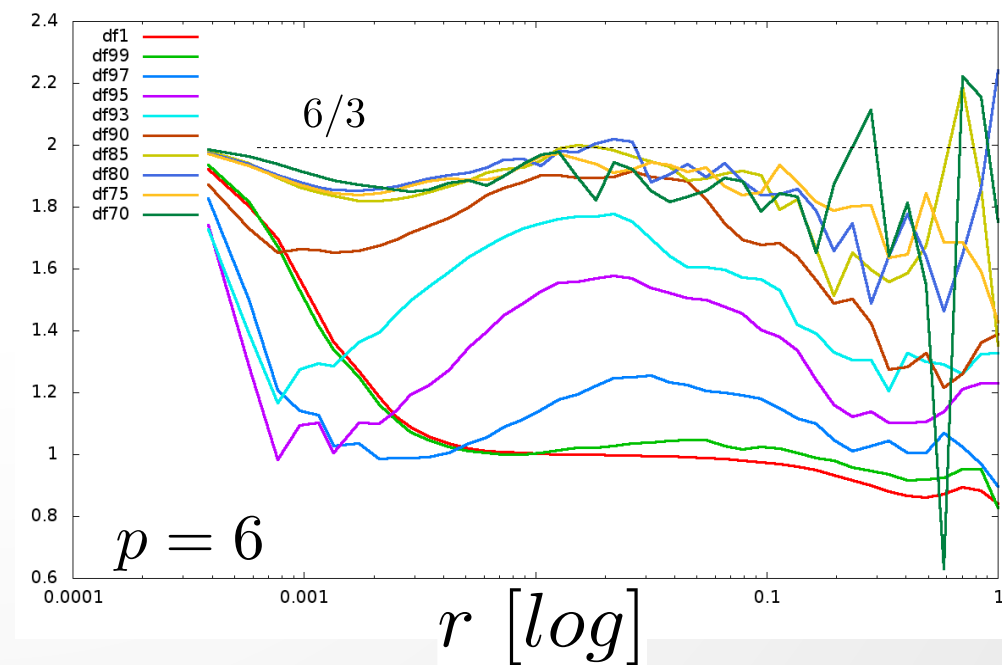
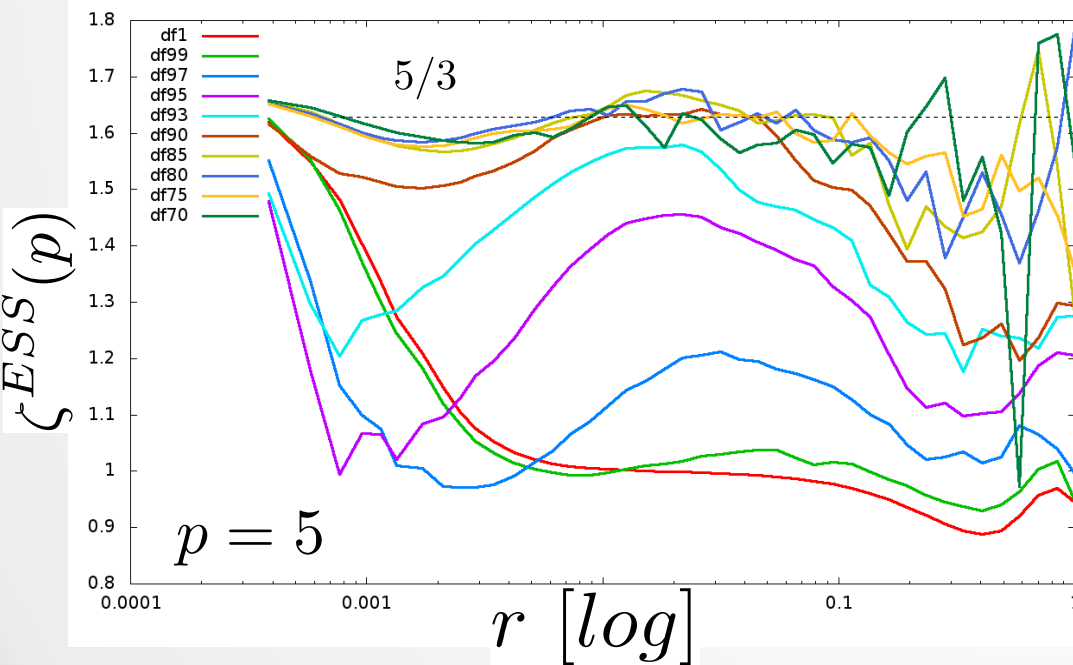
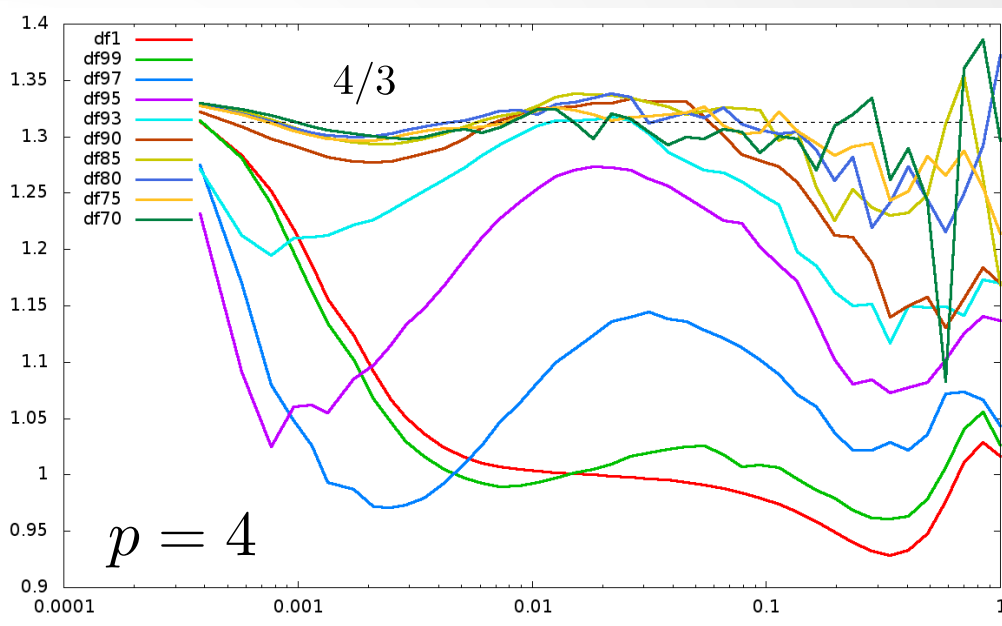
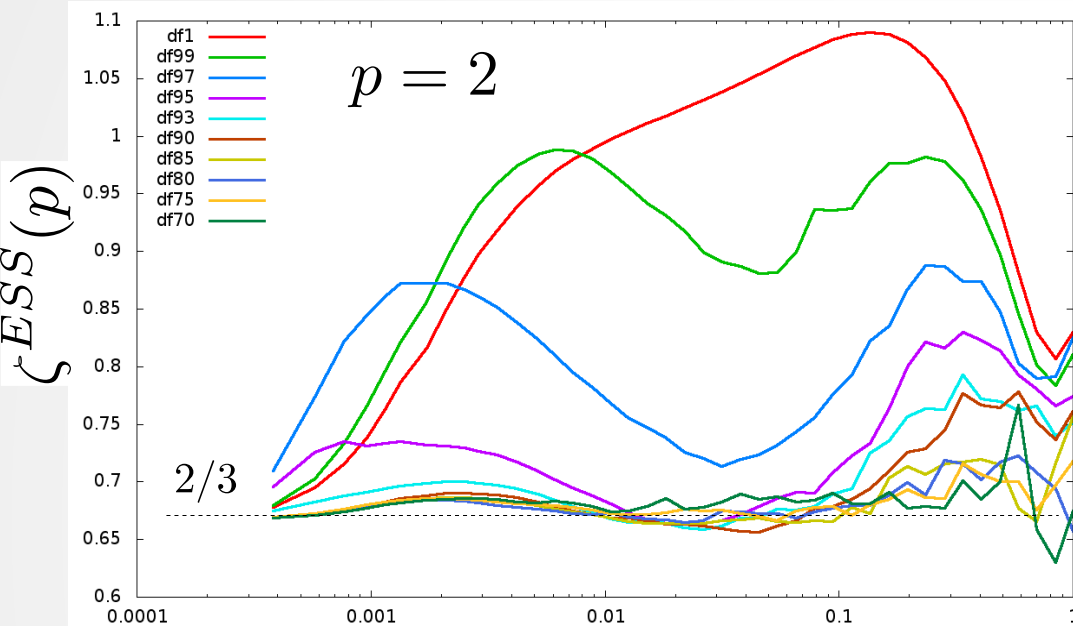


Extended Self Similarity (ESS):

$S_p(r)$ vs $S_3(r)$

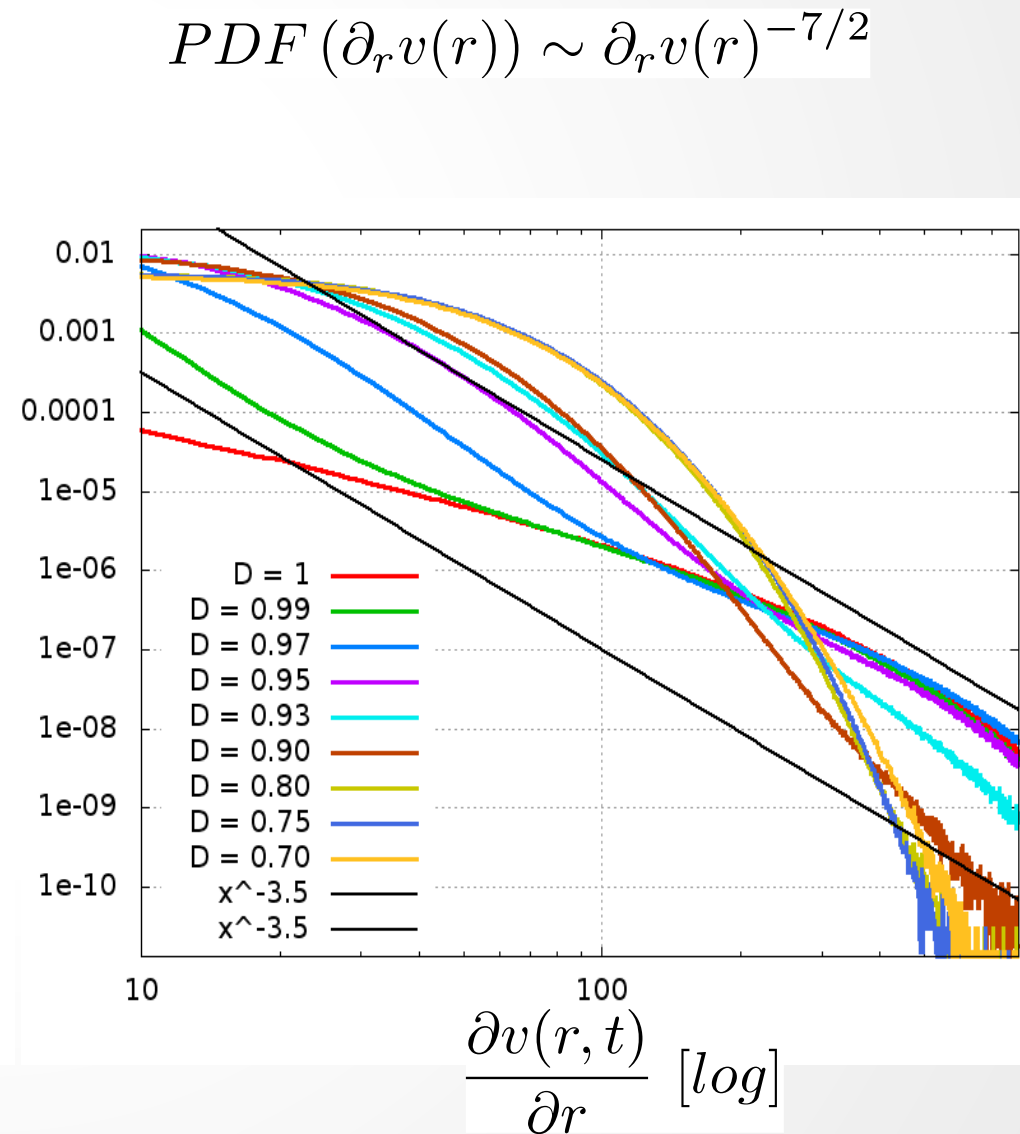
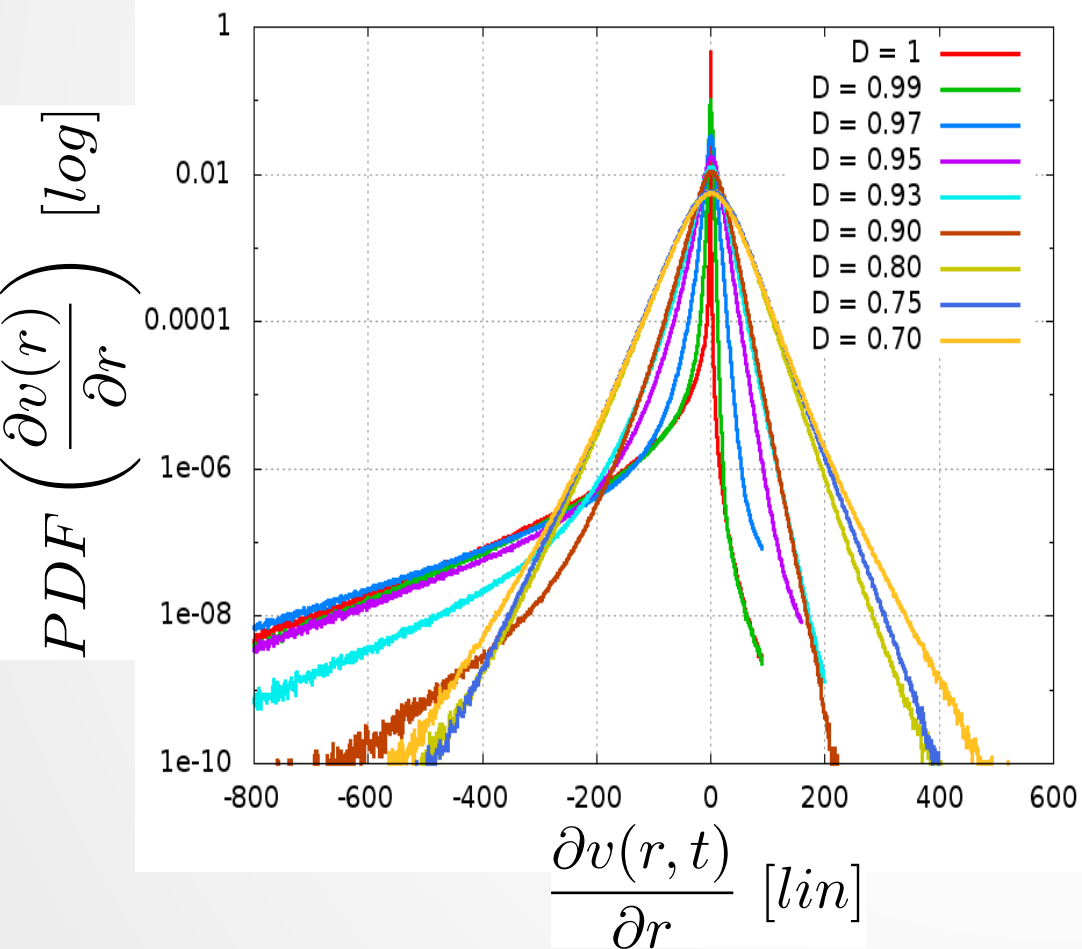
$$\zeta^{ESS}(p) = \frac{\partial \log(S_p(r))}{\partial \log(S_3(r))}$$

$$\zeta^{ESS}(p) = \frac{\zeta(p)}{\zeta(3)}$$



PDF of velocity gradients:

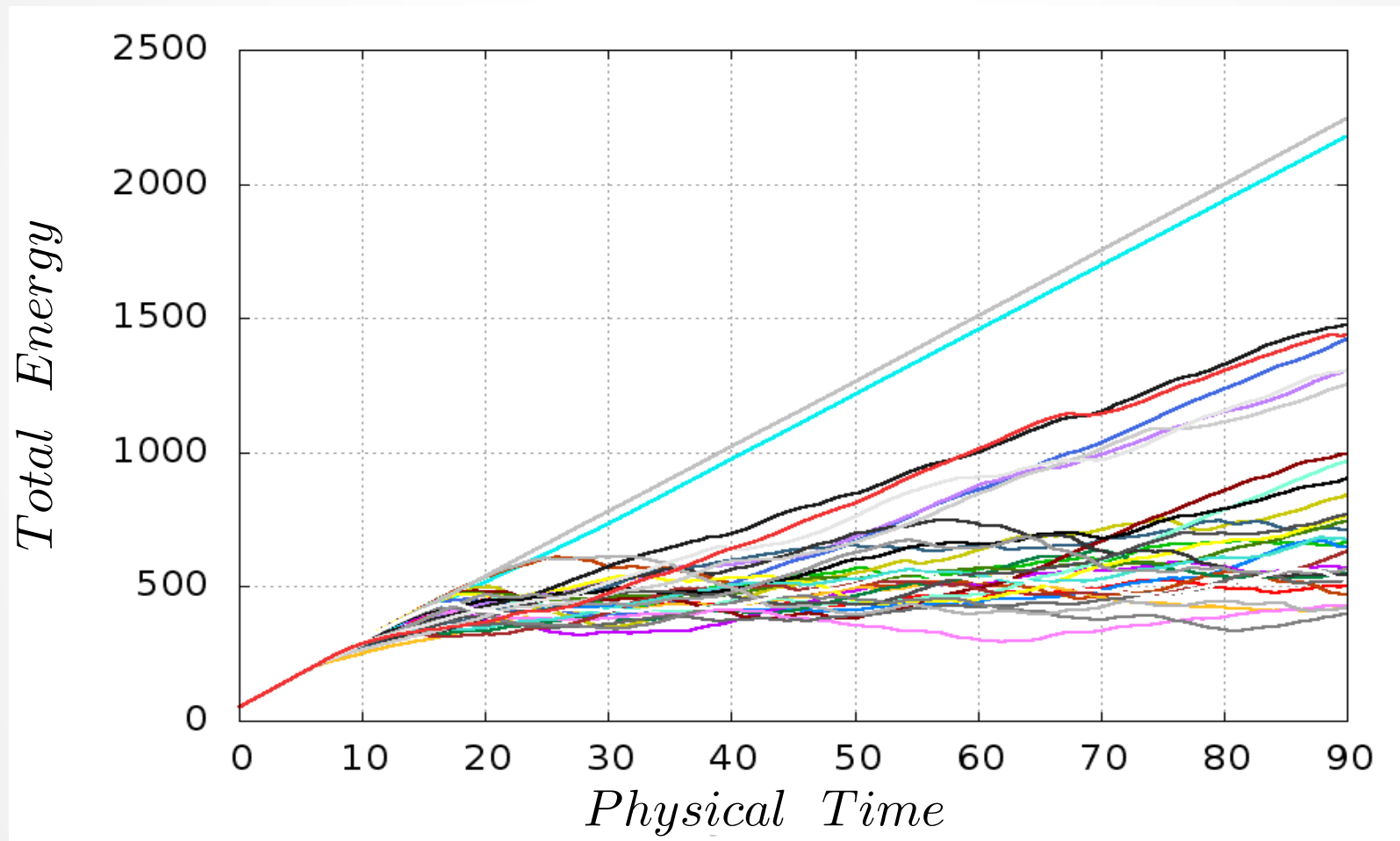
Khanin, Mazel
and Sinai, Probability Distribution Functions
for the Random Forced Burgers Equation,
(1997 Phys. Rev. Lett. 78, 1904)



Non Self-Averaging Problem

Total energy evolution: **different masks**

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2} + f_{deterministic}(x, t)$$

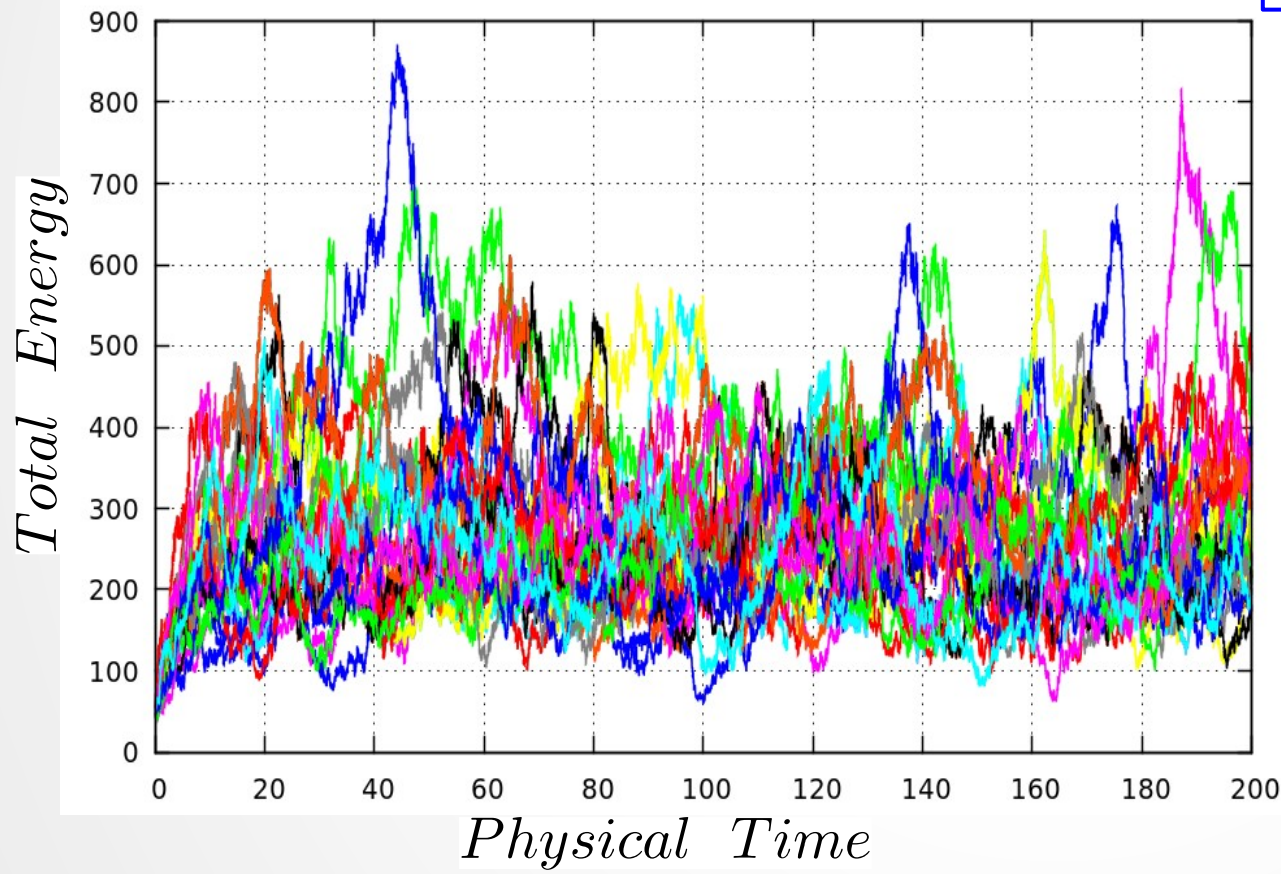


Block in the energy transfer

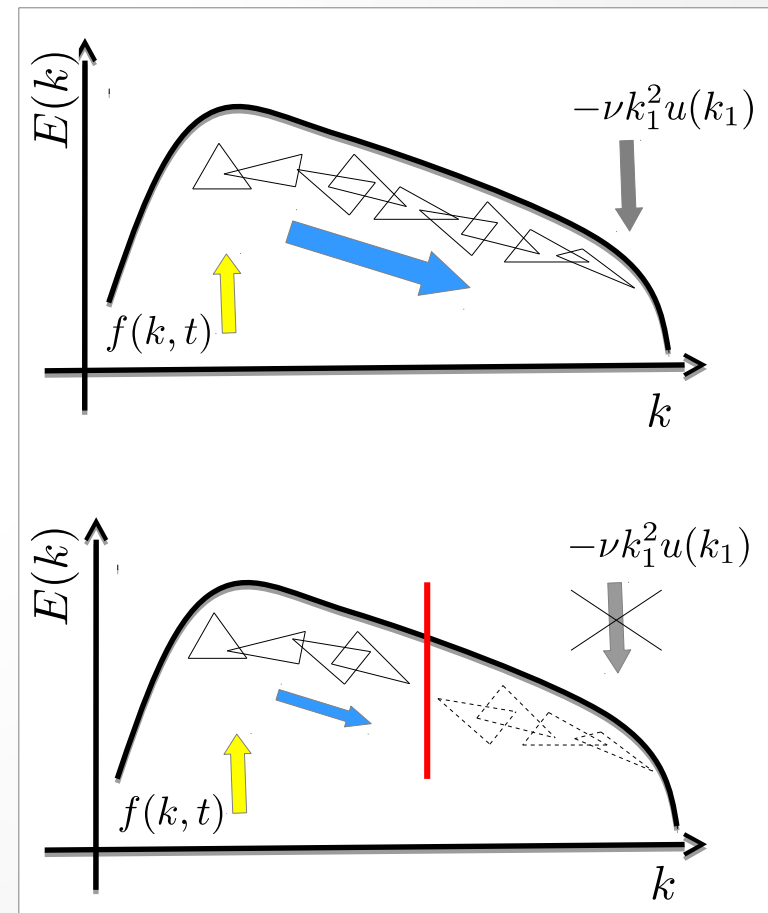
..solution

Total energy evolution: **different masks**

$$\frac{\partial u(k_1, t)}{\partial t} = \sum_{k_2+k_3=k_1} \Pi(u(k_2), u(k_3)) - \nu k_1^2 u(k_1) + f_{stoc}(k, t)$$



Block in the energy transfer



Conclusions:

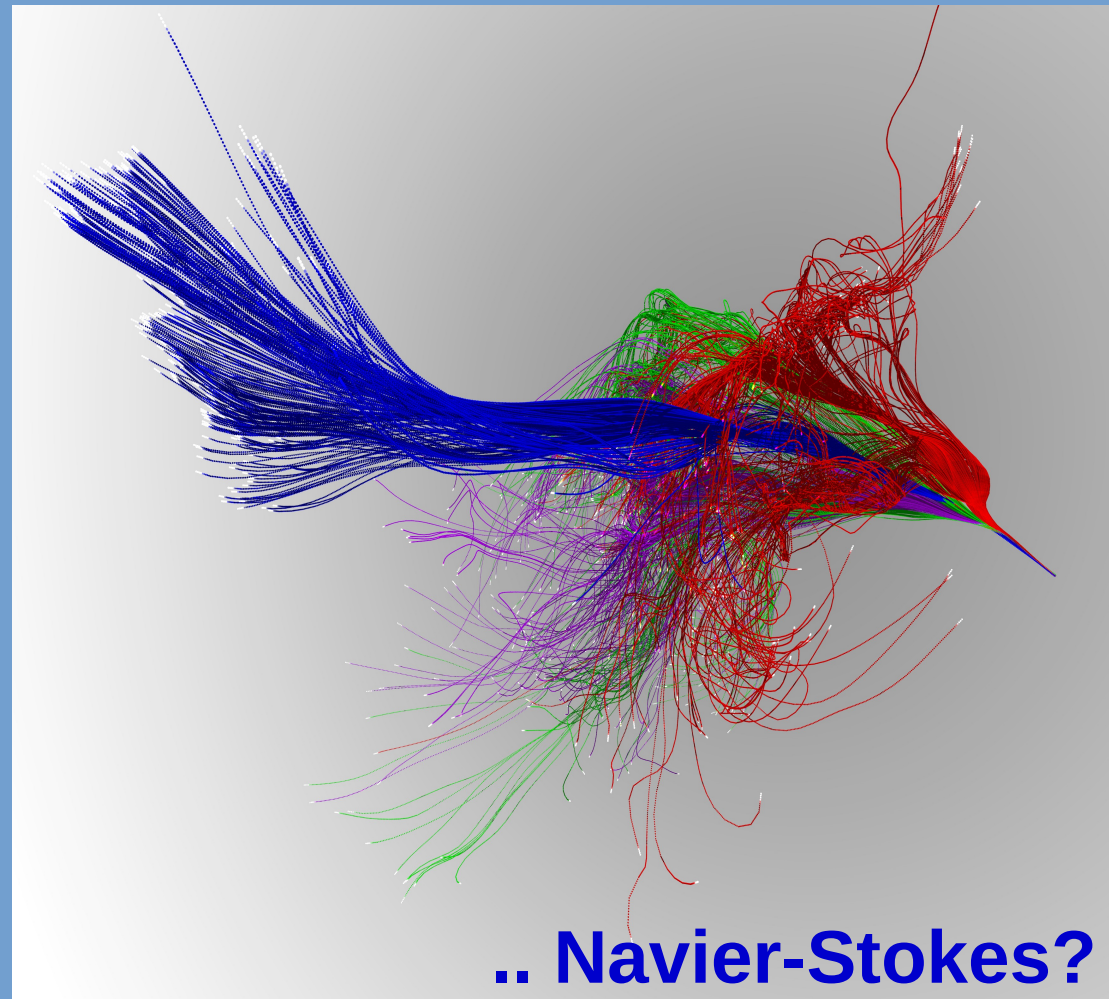
1) Small scales structures appear in the real space velocity field when the system is evolved in a dimension less than 1

2) Differences in the statistic have been found among the different dimensions

- Structure functions
- Pdfs of Gradients

3) The system seems to be no longer intermittent with the decreasing of the dimension starting from a value close to 0.90

4) We need to understand the relation between the scaling laws of 2° order structure functions and energy spectra



Intermittency on Burgers' equation

..BIFRACTAL MODEL

$$\frac{\delta v_\ell(r)}{v_0} \sim \begin{cases} \left(\frac{\ell}{\ell_0}\right)^{h_1}, & r \in \mathcal{S}_1, \dim \mathcal{S}_1 = D_1 \\ \left(\frac{\ell}{\ell_0}\right)^{h_2}, & r \in \mathcal{S}_2, \dim \mathcal{S}_2 = D_2 \end{cases}$$

$$\begin{cases} D_1=0 ; h_1=0 & \leftarrow \text{isolated shock} \\ D_2=1 ; h_2=1 & \leftarrow \text{smooth ramps} \end{cases}$$

$$\frac{\langle \delta v_\ell^p \rangle}{v_0^p} \propto \left(\frac{\ell}{\ell_0}\right)^{ph_1} \left(\frac{\ell}{\ell_0}\right)^{1-D_1} + \left(\frac{\ell}{\ell_0}\right)^{ph_2} \left(\frac{\ell}{\ell_0}\right)^{1-D_2}$$

$$\frac{\langle \delta v_\ell^p \rangle}{v_0^p} \propto \left(\frac{\ell}{\ell_0}\right)^1 + \left(\frac{\ell}{\ell_0}\right)^p$$

Probability to be within a distance l of the set \mathcal{S}

